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THE REAL PART OF THE P-P AND P-D FORWARD SCATTERING AMPLITUDES FROM 50 TO 400 GeV

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ABSTRACT

Proton-proton and proton-deuteron elastic scattering has been measured for incident proton laboratory energy from 50 to 400 GeV. Both reactions were investigated by detection of the recoil particles. The minimum |t| value in the p-p data was 0.0005 (GeV/c) 2 and 0.0008 (GeV/c) 2 for the p-d data; values which lie within the Coulomb region. From these experiments we have determined the ratios of the real part to imaginary part of the forward scattering amplitude, ρ_{pp} and ρ_{pd} , for p-p and p-d scattering. We find that ρ_{pp} and ρ_{pd} show systematic differences that are explicable in terms of a model which assumes $\rho_{pp} = \rho_{pn}$. Using a Glauber approach and a sum of exponentials for the deuteron form factor we obtain ρ_{pn} for p-n scattering.

I. INTRODUCTION

Previous experiments at the Fermi National Accelerator Laboratory have yielded measurements of the ratio of the real part to imaginary part of the forward scattering amplitude, ρ_{pp} , for pp elastic scattering $^{1/2}$ as well as the slope of the forward diffraction peak, b, of p-p³ and p-d elastic scattering. Recent measurements have been reported from the ISR at CERN on ρ_{pp} and the resultant dispersion relation implications on σ_{pp} at ultra high energies. Additional measurements at Serpukhov on p-d elastic scattering below 70 GeV have yielded a parameterization of the deuteron form factor |S(t)|, and the ratio of the real part to imaginary part ρ_{pd} , for p-d scattering. Further analysis using a Glauber approach has resulted in values reported from the Serpukhov data from the ratio of the real part to imaginary part ρ_{pn} , for p-n scattering.

We have extended these measurements to higher energies and with lower momentum transfer, |t|, further into the Coulomb and Coulomb-nuclear interference regions. It is observed that ρ_{pp} rises and crosses through zero at approximately 335 GeV, a consequence and reflection of the rise in the total cross section. Isotopic spin invariance in strong interactions makes the prediction that asymptotically at high energies we would expect to find $\rho_{pp} = \rho_{pn}$. Our motivation was to test this prediction at high energies and in a single experimental setup.

In Section II we describe the experiment and details of the analysis. In Section III we tabulate our p-p data at 52, 80, 199, 261, 303 and 398 GeV and our p-d data at 49, 82, 182, 281, 379 and 397 GeV. In Section IV we discuss the results of the fits and obtain values for the real part of the forward scattering amplitude for p-p and p-d scattering. In this section the deuteron is considered a single target particle. We discuss the energy dependence of ρ_{pp} and ρ_{pd} . In Section V we use the Glauber approach to fit ρ_{pn} . In Section VI we summarize our conclusions.

II. EXPERIMENTAL PROCEDURE AND APPARATUS

Figure 1 is a schematic layout of the experiment. The circulating beam in the main ring of the Fermilab accelerator intercepted a low density gas jet target of hydrogen or deuterium.

Recoil particles were detected at a distance of 7.5 meters by sets of solid state silicon detectors. The detectors, detector box, and ion guide were under vacuum and mounted on a movable assembly; the angle could be varied remotely. The detector, jet profile size, and 7.5 meters distance determine an angular resolution of ±0.8 mrad, a factor of 3 improvement over previous experiments.

The |t |values studied were 0.0005 < |t | < 0.03 (GeV/c) 2 for hydrogen, corresponding to recoil angles 11 < ω < 90 mrad and ranges in silicon 2 < R < 1600 μ m. The |t |values for deuterium were 0.0008 < |t | < 0.08 (GeV/c) 2 corresponding to recoil angles 8 < ω < 75 mrad and ranges in silicon 3 < R < 1500 μ m.

The typical jet operating parameters were density 2 x 10^{-8} g/cm³, jet width (RMS) ± 6 mm, and jet pulse length 100 msec. At this density and at |t| = .0006 (GeV/c)² (pp) or |t| = .001 (GeV/c)² (pd) the correction to do/dt due to multiple scattering of the outgoing recoil particle was about 3%. The Fermilab main circulating beam had typical operating parameters: intensity 1.8 x 10^{13} protons/pulse, beam size 2 mm diameter. Data runs were taken simultaneously at different energy points along the magnetic ramp rise in the range 50 GeV < E_{1ab} < 400 GeV. During the pp runs 3 jets/pulse were used; the pd data was obtained using 4 jets/pulse. The magnetic field in the accelerator was recorded every 30 msec, or every -4 GeV in the incident beam energy.

The recoil protons or deuterons were detected by sandwiches of two solid-state detectors. The detectors used were totally depleted surface barrier silicon with typical dimensions 4 x 20 mm². The front detectors ranged from 15 µm to 250 µm thick and the back detectors from 200 µm to 1500 µm. The recoil particle usually stopped in the back detector. For the lowest | t | range the particle range was less than the thickness of the front detector. The detectors were movable remotely and set at varying angles during the experiment. Two stacks of detectors were permanently fixed and were used to monitor the jet-beam interaction rate.

A schematic diagram of the electronics is shown in Fig. 2. Pulses from the detectors were fed to preamplifiers and then

to amplifiers. The amplified ΔE signal from a given front detector was discriminated and used to trigger a commercial CAMAC peak sensing ADC. During readout of a stack the inputs to all other stacks were inhibited. Thus all channels had the same percentage dead time (< 3%).

The ADC's for the front and back detectors were each supplemented with a TDC which measured the time difference between the front and back detector signals. The scanner scanned the ADC and TDC outputs, and these outputs plus stack identification bits were stored as two CAMAC words in one of two fast 64 word buffer memory units. When one buffer memory unit was full it was disabled and read into a PDP-11 computer. During this time the other buffer memory unit was enabled and data collection continued without interruption. After storage on a magnetic disk, events were written onto a magnetic tape for permanent storage. In addition, the computer generated the gates and timing for the experiment, event histograms, stack display patterns, beam spill histograms, and calibration tasks. Typical data rates were 1,000 events per jet pulse distributed over 8 detector stacks.

The linearity of the electronics was checked by sending test pulses into test inputs of the preamps and recording the output pulse heights. The ADC integral linearity was 0.2%. The differential linearity was 1.5%.

III. DATA ANALYSIS AND CROSS SECTION TABLE

A. Data Analysis

The bit patterns were sorted into 256 x 256 plots of the front detector ΔE vs. the back detector $E-\Delta E$. The mass of a particle stopping in the back element is given by the empirical formula

$$\mathbf{m} = \mathbf{m}_{\mathbf{p}} \left[\frac{\alpha}{d_{\mathbf{F}}} \left[(\mathbf{T}_{\mathbf{F}} + \mathbf{T}_{\mathbf{B}})^{\beta} - \mathbf{T}_{\mathbf{B}} \right] \right]^{1/(\beta-1)}$$
(1)

where α = 0.0133, β = 1.73, d_F is the thickness of the front detector in mm, and $T_F(T_B)$ is the energy deposited in the front (back) detector in MeV. Using Eq. (1) we define deuterons by the cuts 1.5 < m/m_p < 2.5 and protons by the cuts 0.5 < m/mp < 1.5.

The momentum distribution of the recoil is shown in Fig. 3. The momentum spectra were fitted over the range > 5.0 σ (σ is the Gaussian standard deviation) to a formula which contained Gaussian plus exponential (background) terms. The number of elastic events was obtained after applying cuts at ±4 σ and subtracting the background, which was determined from the fit. The background was ~1% except for the lowest | t | deuteron data where background was ~3%. We have checked that fits made with different "smooth" background terms (linear, quadratic) give approximately the same result. Our low background is due to the long ion guide and good shielding. The width of the elastic

peak is consistent with a total angular resolution of ± 0.8 mrad expected from our jet size (± 6 mm) and detector geometry (± 4 mm). The good angular resolution led to a negligible (< 0.28) inelastic background under the elastic peak, arising from reactions pp + Xp or pd + Xd.

The number of elastic events in each detector, ΔN , in solid angle $\Delta \Omega$, were converted to $d\sigma/dt$ as follows:

$$\frac{d\sigma}{dt} = \frac{\Delta N}{\Delta \Omega} \frac{\pi}{2m \sqrt{|t|}} \left(\frac{d\sigma}{d\Omega}\right)_{\text{fixed}} / \left(\frac{\Delta N}{\Delta \Omega}\right)_{\text{fixed}}$$
 (2)

where $\left(\frac{d\sigma}{d\Omega}\right)_{\text{fixed}}$ and $\left(\frac{\Delta N}{\Delta\Omega}\right)_{\text{fixed}}$ are for the fixed monitors. $\frac{\Delta N}{\Delta\Omega}$ and $\left(\frac{\Delta N}{\Delta\Omega}\right)_{\text{fixed}}$ are the number of elastic events after applying the geometrical and resolution correctons. At the smallest |t| values our geometrical correction was about 1.5% and correction due to resolution about 7%. The correct normalization was obtained in the final fits from the optical theorem.

B. Systematic Errors

1. Angle

The absolute angle of the detector assembly was determined from the energy of the elastic peaks wholly contained in the front detectors. Detectors were calibrated with an alpha particle source, 90Th (5.446, 5.684, 6.04, 6.277, 6.774, 8.786 MeV). The absolute angles determined from the elastic peaks and alpha particle energy calibrations, when compared with angle

data from survey measurements show our estimated angular uncertainty to be ≤ 0.15 mrad. Recent measurements with p/d and p/He target mixtures indicate that this angle offset is independent of the particle type, range, or detector angle. Shielding reduced the magnetic field in our detector system to ≤ 0.03 gauss. This minimized angular errors at very low |t|. Due to our angular uncertainty the systematic error on $d\sigma/dt$ in the interference region is about 1%.

2. Area

The uncertainty in detector area is ±0.15%.

3. Nuclear Interactions

The effect of deuteron and proton interactions in the silicon detector are approximately proportional to the particle range. For our t range the effects are ≤ 0.07 %.

C. Cross Section Table

The cross sections for p-p at 52, 80, 199, 261, 303, 398 GeV and for p-d at 49, 82, 182, 281, 379, 397 GeV are given in Table I. The errors listed are statistical. The absolute normalizations, determined in the fits from the optical theorem point, are uncertain by =0.7% for p-p and =0.5% for p-d scattering. The p-p data at 398 GeV and the p-d data at 379 GeV are shown in figures 4 and 5.

IV. DETERMINATION OF ρ_{pp} AND ρ_{pd}

Our results for p-p elastic cross sections, listed in Table 1, have been fitted to the Bethe interference formula which includes electromagnetic effects.

$$\frac{d\sigma}{dt} = \pi |f_h + f_c|^2$$
 (3)

where the nuclear and Coulomb scattering amplitudes take the forms

$$f_n = \frac{\sigma_{\text{tot}}}{4\pi N} (\rho_{pp} + i) e^{bt/2}$$
 (4)

$$f_{c} = \frac{2\alpha i}{t} G_{p}^{2}(t) e^{i\alpha j}$$
 (5)

The free parameters in the fit are ρ_{pp} and the overall normalization which sets a scale to the values listed in Table I. $\Delta\rho/\Delta b$ and $\Delta\rho/\Delta\sigma$ can be used to calculate the change in ρ due to a change in the assumed values of b and σ_{tot} .

In Eq. (4) we have assumed that the real and imaginary parts have the same t dependence and that spin effects may be neglected. For $\sigma_{\rm tot}$ and b we use empirical fits to the data of Carroll et al., $\sigma_{\rm tot}$ (mb) = 50.866 - 5.2302 ln Spp + 0.5437 ln Spp, and Bartenev et al., b(s) = 8.27 + 0.556 ln Spp. In Eq. (5) $\sigma_{\rm tot}$ is the fine structure constant $\sigma_{\rm p}$ (t) = $\sigma_{\rm tot}$ (1 + |t|/0.71) $\sigma_{\rm tot}$

the proton electromagnetic form factor, and $\alpha \phi = 2\alpha \ln \frac{1.06\text{M}}{\text{R}^3/\text{t}}$ the Coulomb phase with $R = \sqrt{10} \text{ mb}^{1/2}$.

The χ^2 minimization results for ρ_{pp} the ratio of the real to imaginary parts of the forward scattering amplitude, are listed in Table II. Typical χ^2 's are 1.2-1.3 per degree of freedom. In Fig. 6 we show these values together with previously published results. Results on ρ_{pp} from the Fermilab experiment are not shown on Fig. 6 since too many points in the same energy region would be confusing. Our results are in agreement with the previous experiment and have somewhat smaller statistical and systematic errors. Furthermore in the present case tables of differential cross sections are given so readers may fit other models to the data. An empirical expression valid within our energy range is

$$\rho_{pp}(s) = (-0.490 \pm 0.034) + (0.076 \pm 0.006) \ln s_{pp}.$$
 (6)

The solid curve shown in Fig. 6 is derived from a dispersion relation

Re A = C +
$$\frac{E}{4\pi^2}$$
 $\int_{m}^{\infty} dE' \left[\frac{\overline{\sigma}_T(E')}{E' - E} - \frac{\overline{\sigma}_T(E')}{E' + E} \right]$ (7)

where σ_T and $\bar{\sigma}_T$ are proton-proton and anti-proton-proton total cross sections. The normalization of the curve (adjustment of C) is arbitrary using our data and that of Amaldi ct al. The dashed curve in Fig. 6 is calculated assuming only vacuum exchange.

Our results for p-d elastic cross sections, listed also in Table I, have been fitted with Eq. (3), the Bethe interference formula, where we treat the deuteron as a single particle. In this case the nuclear amplitude is parameterized as a sum of exponentials*,12

$$f_{n} = \frac{\sigma_{tot}}{4\pi \hbar} (\rho_{pd} + i) e^{bt/2}$$

$$x = \frac{141.5t}{0.34e} + 0.58e \frac{26.1t}{4} + 0.08e$$
(8)

where we assume that b(s) = b_{pp} = 8.27 + 0.556 ln S_{pp} and use an empirical fit to the published pd total cross section^{10,13} $\sigma_{tot,pd}$ (mb) = 99.73 - 9.40 ln S_{pd} + 0.829 ln $^2S_{pd}$. The Coulomb amplitude is written as

$$f_{c} = \frac{2\alpha H}{t} G_{p}(t) G_{d}(t) e^{i\alpha \beta}. \tag{9}$$

with
$$G_d(t) = e^{(25.95t + 60t^2)/2}$$
 and $\alpha \beta = 2\alpha \ln \frac{1.06 \text{M}}{\text{RV}/|t|}$ with $R = 2.7 \sqrt{10} \text{ mb}^{1/2}$.

In Table II we list our results for the ratio of the real to imaginary parts of the forward scattering amplitude, ρ_{pd} . An empirical expression valid in our energy range is:

$$\rho_{pd} = (-0.450 \pm 0.035) + (0.070 \pm 0.006) \text{ fin } S_{pp}$$
 (10)

The main contribution to the systematic error in ρ_{pp} (±0.009)

and in ρ_{pd} (±0.008) is the angular uncertainty (0.15 mrad) in the position of the detectors. In Fig. 7 we show our ρ_{pd} values together with previously published results. The solid line shown is our ρ_{pp} (s) parameterization, Eq. (6).

Pertinent to our understanding of the relationship between $\rho_{pd} \text{ and } \rho_{pp} \text{ is the following Glauber derivation. If one expresses }$ the deuteron forward scattering amplitude as

$$f_d = f_p + f_n + i / (r^{-2})_d f_p f_n$$
 (11)

where we know the deuteron inverse radius squared, $< r^{-2} >_d = I_G = 0.033 \text{ mb}^{-1}$. ($I_G \text{ will be discussed in the next section.})$

If we further assume that the forward proton and neutron amplitudes are equal

$$f_n = f_p = \frac{\sigma_{tot}}{4\pi k} (\rho_{pp} + i)$$
 (12)

then substitution of Eq. (12) into Eq. (11) and omission of the higher order ρ_{pp}^{-2} term, will give

$$f_{d} = \frac{\sigma_{tot}}{2\pi N} \left[\rho_{pp} (1 - \delta) + i (1 - \frac{\delta}{2}) \right].$$
 (13)

In the above we use the notation of a Glauber screening parameter $\delta \equiv < r^{-2} > \frac{\sigma_{tot}}{4\pi} \simeq 0.111.$ The implication of Eq. (13) is that we would expect

$$\rho_{pd} = \frac{\text{Re } f_d}{\text{Im } f_d} = \rho_{pp} - \frac{(1 - \delta)}{(1 - \frac{\delta}{2})} = 0.941 \rho_{pp}. \tag{14}$$

This is a relationship satisfied by the data in Table II and Table III and displayed in the lines drawn on Fig. 7. Both ρ_{pp} and ρ_{pd} cross zero at 335 GeV. Equation (14) predicts that ρ_{pd} and ρ_{pp} will be zero at the same energy. Using the empirical fits to our data ($\rho = \rho_{o} + \rho_{1}$ in s) as expressed by Eqs. (6) and (10) we find

$$\rho_{\rm 1pd}/\rho_{\rm 1pp} = 0.92 \pm 0.11$$
 (15)

in good agreement with the prediction. Since the derivation is based upon the equivalency of the proton and neutron our results support the idea of isotopic spin invariance in proton neutron collisions at high energies.

V. THE GLAUBER APPROACH AND $\rho_{\mathbf{p}n}$

In the Glauber approach 12,14-16 elastic p-d scattering is described as a coherent sum of Coulomb, single-nucleon, and double-nucleon scattering. Assuming only s-wave contribution to the form factor and only elastic rescattering we can write

$$\frac{d\sigma}{dt} = \pi |S(t/4)| (A_C + A_D + A_D) + A_G^{-1/2}, \qquad (16)$$

where

$$A_{C} = \frac{2\alpha N}{t} e^{b}pp^{t/2} e^{i\eta}$$

$$A_{p} = e^{i\chi}cp \frac{\sigma_{pp}}{4\pi N} (\rho_{pp} + i) e^{b}pp^{t/2},$$

$$A_{n} = e^{i\chi}cn \frac{\sigma_{pn}}{4\pi N} (\rho_{pn} + i) e^{b}pn^{t/2},$$

$$A_{G} = e^{i\chi}cpn \frac{iN \sigma_{pp}\sigma_{pn}}{(4\pi N)^{2}} (\rho_{pp} + i) (\rho_{pn} + i) e^{t(b}pp + b_{pn})/8 I_{G}.$$

S(t/4) is the deuteron form factor, b_{pp} , b_{pn} , ρ_{pp} , ρ_{pn} and σ_{pp} , σ_{pn} are the slope parameters, real to imaginary forward nuclear scattering-emplitude ratios, and total cross sections for pp and pn scattering, respectively, and η , χ_{cp} , χ_{cn} , and χ_{cpn} are phases between the amplitudes. We assume t independence for ρ_{pp} and ρ_{pn} , as no experimental information exists on this point. I_{G} is the Glauber intergral:

$$I_G = \frac{1}{2} \int_{\infty}^{0} S(t) e^b p p^{t/2} e^b p n^{t/2} dt.$$
 (17)

We calculate phases using the formulas in Ref. 12:

$$x_{cp} = 0.10, x_{cn} = 0.11, x_{cpn} = 0.10,$$

$$\eta = (2/137.03) \left\{ \ln 2p_{lab} / t \right\} - 0.577$$

For the p-p total cross section we refer to the empirical formula given previously. We assume that $\sigma_{pn} = \sigma_{pp}$. For b_{pn}

= b_{pp} we take the empirical fit, b_{pp} = 8.27 + 0.556 $ln s_{pp}$. For ρ_{pp} we use our measured values summarized in the empirical fit, Eq. (6). The spherical deuteron form factor, s_{o} (t/4), is given by a sum of exponentials

$$S(t/4) = S_0 (t/4) = A_1 e^{B_1 t/4} + A_2 e^{B_2 t/4} + (1-A_1-A_2) e^{B_3 t/4}$$
 (18)

with the coefficients previously given in Eq. (8). Calculations show that inclusion of a quadrupole form factor is significant only at large |t| values and introduces a maximum contribution of ≈ 1.68 at our highest |t|.

The results of the Glauber analysis are given in Table III and shown in Figs. 8 and 9. The free parameters are ρ_{pn} and I_G . One notes a larger statistical error on ρ_{pn} in comparison with the ρ_{pd} statistical error. This is inherent in the analysis. Not shown in Fig. 8 are the systematic errors in our ρ_{pn} values introduced by uncertainty in the deuteron form factor S(t/4). These systematic errors are of the order of our statistical errors. In Fig. 8 is shown the total error corridor on our ρ_{pp} results for comparison of the ρ_{pn} values with ρ_{pp} . The shadow correction is shown in Fig. 9. We have also tried other hypotheses for b_{pn} , e.g. using the empirical fit found in Ref. 4, b_{pn} = 5.57 \pm 1.15 \ln S_{pp} . Using this shrinkage instead of the more reasonable b_{pn} = b_{pp} shifts the ρ_{pn} values up by about one standard deviation.

In (16) we neglected the Coulomb amplitude contribution to the double scattering term. This problem was considered by Franco and Varma. For pd elastic scattering in the Serpukhov energy range they estimated that the difference between the Glauber formula which includes all double scattering terms and the more approximate Eq. (16) is less than 2% for d σ /dt in the Coulomb-nuclear interference region. The difference in the value of $\rho_{\rm pn}$ obtained from the two formulas is ~0.03. Another complication is that in our energy range inelastic intermediate states may contribute to the double scattering term. It is not expected that inelastic and elastic contributions have the same t-dependence.

To investigate how the problems mentioned above may affect our result on ρ_{pn} we also determined ρ_{pn} by another method using the Glauber formula for the pd elastic scattering forward amplitude. Using (11) and the optical theorem one can derive the following:

$$\rho_{pn} = \frac{1}{\sigma_{pn}} \left\{ (\sigma_{pd} \rho_{pd} - \sigma_{pp} \rho_{pp}) \left(1 + \frac{\langle r^{-2} \rangle_{d}}{4\pi} \sigma_{pp}) + (\sigma_{pd} \sigma_{pp}) \frac{\langle r^{-2} \rangle_{d}}{4\pi} \sigma_{pp} \rho_{pp} \right\}$$

$$(19)$$

To estimate ρ_{pn} we used ρ_{pd} given in Table II and parameterized σ_{pp} , σ_{pd} and ρ_{pp} as described in Section IV. We assume $\sigma_{pn} = \sigma_{pp}$ and values for < r $^{-2}$ >_d = I_G were taken from Table III. The systematic error in ρ_{pn} due to the uncertainty in the double

scattering term is energy dependent and small compared to the statistical errors. If we assign an error of ~20% to I_G then the resulting systematic error in ρ_{pn} is < .005. Comparison of ρ_{pn} obtained from (16) (Table III) with those obtained from (19) (Table IV) show no significant difference between the two methods.

VI. CONCLUSIONS

Comparison of our pp data with dispersion relation calculations show good agreement. By comparing our pp data with the prediction obtained from a simple vacuum exchange model we can estimate the importance of other exchanges (Fig. 6). ρ_{pp} and ρ_{pd} both cross zero at about 335 GeV. In Section V we showed that using a Glauber approach to obtain ρ_{pn} we found no significant difference between ρ_{pp} and ρ_{pn} . The shadow correction increases with energy (Fig. 9) but remains small, \leq 0.03 mb⁻¹.

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the fit then ρ_{pd} = (-0.473 ± .052) + (0.073 ± 0.007) £n S_{pp}.

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	[(GeY/c)']	du/dt] [mb/(GeV/c) ²]	/(da/dt) [mb/(GeV/c)*]		-t [(GeV/c) [*]]	da/dt [mb/(GeV/c) ⁷	/(do/dt)] [mb/(GeV/c) ²]
	0.00063 0.00065	765.90	17.34		0.00077	756,23	21.98
	0.00076	724.49 553.59	16.61		0.00100	596.68	14.61
	0.00100	366.71	7.54 5.50		0.00150	395.00	8.97
	0.00124	263.87	5.4}		0.00178	341.63	13.20
	0.00131	257.86	7.28		0.00184	334.00	17.17
	0.00132	244.52	3.99		0.00212	360.62	12.88
PP + PP	0.00133	248.51	7.11	-44	0.00219	313.55	12.26
	0.00159	203.57	5.66	pd → pd	0.00235	311.44	5.60
i2 GeV	0.00161	191.74	5.44	49 GeV	0.00243	320, 19	5.68
	0.00162	191.11	3.39	43 GEY	0.00268	299.16	4.58
	0.00193	161.20	3.82		0.00284 0.00291	289.46	8.86
rmalization		152.86	2.38	flormal ization	0.00291	293.72	8.88
certainty	0.00214	149.05	3.25	uncertainty	0.00357	286.02	4.31
3 000	0.00223	150.19	5.25		0.00407	281.07	7.68
0.90%	0.00225	143.14	5.12	0.58%	0.00412	258.64 253.59	7.54
	0.00242	129.94	2.66	0. 50%	0.00415	259.15	4.42
	0.00248	136.74	3.88		0.00438	254.69	7.55
	0.00249 0.00280	134.87	2.78		0.00456	249.00	6.40
	0.00302	117.62	2.58		0.00502	243.61	3.67
	0.00302	115.44	2.80		0.00510	239.98	5.79 4.32
	0.00339	101.05 101.73	3.19		0.00588	223.14	3. 6 8
	0.00344	108.44	3.19		0.00671	219.15	3.81
	0.00383	98.71	2.47		0.00749	197.37	5.83
	0.00386	102.74	3.14		0.00929	206.45	8.39
	0.00409	94.06	3.21 2.07		0.01070	178.40	3.26
	0.00435	102.48	2.72		0.01152	176.33	5.28
	0.00457	92.73	1.90		0.01227	174.95	6.29
	0.00466	94.36	1.98		0.01402	156.46	3.49
	0.00483	89.36	3.09		0.01484	136.61	7.13
	0.00486	90.75	3.12		0.01500	150.11	7.64
	0.00518	88.12	1.90		0.01529	150.54	3.13
,	0.00561	90.78	2.15		0.01563	146.17	3.38
	0.00593	83.83	1.91		0. 01662	139.45	2.77
	0.00633	87.52	2.00		0.01766	138.27	4.00
	0.00647	82.55	2.88		0.01943	126.05	2.66
	0.00651	85.63	1.96		0.02056	125.10	4.12
	0.00652	84.86	2.93		0.02126	116.61	2.03
	0.00708	81.54	2.23		0.02264	106.90	3.11
	0.00713	82.76	2.25		0.02383	107.58	2.30
	0.00739	80.21	1.61	•	0.02490	102.38	3.89
	0.00780	82.02	2.01		0.02511 0.02610	102.29	3.91
	0.00803	77.13	1.58		0.02851	99.90	1.77
	0.00842	77.00	2.5 9		0.02874	87.41	3.34
	0.00889	78.85	1.48		0.02914	91.93	3.32
	0.00944	77.95	2.25		0.02961	86.07 84.86	1.88
	0.00946	77.32	1.69		0.03096	83.26	1.67
	0.00986 0.01033	77.61	1.93		0.03304	75.81	2.00 1.59
	0.01057	77.87 74.22	1.75		0.03477	69.44	1.59
	0.01061	74.22 75.22	2.26		0.03623	71.12	1.61
	0.01171	74.24	1.51		0.03628	73.56	2.96
	0.01218	73.68	1,65 1.78		0.03653	73.40	3.03
	0.01253	73.66 70.54	1./5		0.03902	62.88	1.37
	0.01340	70.54 73.47	1.28 1.56		0.04137	55.51	1.32
	0.01424	68.36	1.49		0.04192	58.53	1.47
	0.01474	70.66	1.58		0.04197	61.04	2.34
	0.01530	70.41	1.47		0.04225	57.41	2.33
	0.01555	69.13	1.33		0.04744	48.33	1.08
	0.01565	68.91	1.65		0.04803	46.73	1.16
	0.01700	68.05	1.38		0.04808	50.45	2.19
	0.01754	63.63	1.45		0.05392	37.77	0.96
	0.01797	64.95	1.40		0.05456	36.91	0.89
	0.02000	62.77	1.28		0.05482	36.01	1.31
	0.02068	65.33	1.41		0.05487	36.85	1.30
	0.02175	63.97	1.47		0.06151	31.12	0.83
	0.02334	60.02	1.19				
	0.02514	58.32	1.00				
	0.02753	57.49	1.22				
1	0.02877 D.03060	58.32	1.25				
		55.21	1.07				

	0.00066	672.58	13.61		0.00000	C22 00	
	0.00067	655.77	13.43		0.00092 0.00103	632.28 549.71	13.34 12.75
	0.00080 0.00104	494.86	8.50		0.00106	528.00	12.34
	0.00135	344.26 223.31	6.46 5.53		0.00155	402.41	9.58
	0.00137	222.65	3.51		0.00159 0.00207	400.19 332.66	9.46 5.11
pp + pp	0.00164	177.93	4.62		0.00233	318.22	5.44
pp v pp	0.00166 0.00168	184.49 173.73	4.72 3.89	pd → pd	0.00244	312.87	5.71
80 GeV	0.00200	144.74	3.05	82 GeV	0.00255 0.00256	309.11 306.69	4.94 5.46
	0.00211 0.00220	138.91	2.72		0.00261	308.89	5.47
Normalization		137.74 130.45	2.60 2.67	Normalia sada	0.00272	309.04	5.12
unce rtainty	0.00250	124.13	2.58	Normalization uncertainty	0.00354 0.00375	277.05 2 62 .49	4.59 5.13
A 600	0.00288	110.12	2.13	•	0.00382	263.09	5.14
0.69%	0.00311 0.00348	105.17 101.52	2.22 2.59	0.47%	0.00394	253.62	4.13
	0.00351	99.61	2.56		0.00430 0.00445	241.01 247.89	4.42 4.83
	0.00353	97.61	2.06		0.00460	238.91	4.09
	0.0 0392 0.0 0395	96.98 95.84	2.22 2.20	•	0.00483	236.41	4.24
	0.00419	89.58	1.69	•	0.00533 0.00590	240.4 <u>2</u> 221.88	5. 6 6 3.96
	0.00446 0.00468	89.42 88.13	1.44		0.00653	215.14	3.10
	0.00478	86.57	1.64 1.55		0.00691 0.00699	209.24	4.53
	0.00495	88.47	2.42		0.00700	210.84 210.94	3.39 4.56
	0.00498 0.00575	88.69 83.00	2.42		0.00718	209.39	3.62
	0.00648	80.75	1.20 1.52		0.00737	201.09	3.06
	0.00663	81.45	2.16		0.00766 0.00775	201.05 209.04	3.20 2.7 6
	0.00666 0.00667	80.61	1.38		0.00858	199.06	3.26
	0.000724	80.76 76.19	2.15 1.84		0.00899	192.28	3.05
	0.00729	76.80	1.84		0.00967 0.00974	185.44 185.35	3.89 2.71
	0.00755 0.00797	76.16	1.29		0.00977	182.99	3.85
	0.00757	77.09 72.01	1.54 1.29		0.01067	172.35	2.67
	0.01007	72.47	1.42		0.01125 0.01188	175.61 168.33	2.11 4.64
	0.01083	72.07	1.28		0.0 1199	171.87	4.70
	0.01197 0.01243	70.67 70.58	1.24 1.45		0.01208	168.79	3.01
	0.01279	67.51	1.09		0.01363 0.01403	154.66 154.10	2.05 2.33
	0.01294 0.01452	71.37	1.41		0.01445	150.73	2.76
	0.01503	67.81 68.37	1.26 1.44		0.01454	148.29	2.48
	0.01594	63.07	1.30	•	0.01458 0.01638	148.25 138.83	2.73 2.26
	0.01732 0.01788	64.81 65.23	1.24 1.29		0.01667	1.37.69	2.73
	0.01831	64.74	1.29		0.01713 0.01727	137.48	2.48
	0.01858	63.08	1.18		0.01843	137.25 127.24	2.48 2.11
	0.01960 0.02038	63.72 61.29	1.16 1.10		0.01923	124.01	1.90
	0.02106	60.18	1.15		0.01970 0.02134	123.88 121.02	1.67
	0.02215	59.32	1.16		0,02145	112.81	2.66 2.47
	0.02377	58.02 58.74	1.01 1.14		0.02148	116.94	2.60
	0.02802	56.46	1.04		0.02366 0.02457	103.69 100.60	1.59
	0,02928	56.34	1.43		0.02473	101.70	1.97 1.98
					0.02706	90.98	1.83
					0.02763 0.02835	95.69 88.48	1.56 1.31
					0.03129	82.09	1.17
					0.03193	78.34	1.85
					0.03211 0.03369	78.48 75.40	1.85 1.36
					0.03476	71.00	0.98
					0.03758 0.03773	63.78	1.60
					0.03778	65.05 64.67	1.03 1.62
			•		0.03954	63.31	1.07
					0.04066 0.04347	57.93	0.99
					0.04347	52.59 53.96	1.39 1.41
					0.04509	50.35	0.75
					0.04581 0.04674	51.40 47.21	0.91
					0.05355	37.85	0.91 0.78
					0.05489	36.83	0.66
					0.05921 0.06054	31.88 30.03	0.60 0.66
					0.06226	29,21	0.56
					0,06655	25.35	0.57
•			•		0.06797 0.06978	24.49 22.08	0.54 0.46
					0.07583	18.81	0.50
					•		

	0.00068 0.000104 0.00104 0.00129 0.00136 0.00137 0.00138 0.00165 0.00167 0.00169 0.00212 0.00222 0.00231 0.00233 0.00251 0.00258 0.00290 0.00313 0.00356 0.00356 0.00356 0.00356 0.00356 0.00357 0.00580 0.00472 0.00482 0.00499 0.00503	648.06 477.58 319.27 231.56 210.40 215.02 164.55 165.03 132.94 130.00 127.61 119.68 116.37 106.40 107.35 92.59 94.06 98.79 98.79 98.79 98.79 98.70 81.21 77.08 81.21 81.21 77.08 81.21 8	14.79 6.36 4.76 4.630 3.45 6.26 4.76 2.99 3.34 2.81 4.51 4.34 1.99 2.87 1.76 1.83 2.167 2.18 1.74 2.18 1.74 2.18 1.79 1.79 1.79 1.79 1.79 1.79 1.79 1.79	0.00109 0.00159 0.00164 0.00212 0.00239 0.00263 182 GeV 0.00269 0.00269 0.00269 0.00364 0.00364 0.00364 0.00364 0.00365 0.00365 0.00366 0.00366 0.00367 0.00473 0.00473 0.00473 0.00514 0.00514 0.00514 0.00514 0.00719 0.00719 0.00719 0.00719 0.00738 0.00787 0.00787 10.00796 0.00881 0.00924 0.00924 0.00924 0.00924 0.00981 0.0093 0.01003 0.01003 0.01003 0.01003 0.01003 10.01156 0.01220 0.01231 0.01156 0.01220 0.01440 0.01497 0.01497 0.01497 1.001682 0.01493 1.001759 1.30 0.01497 1.50 0.01497 1.50 0.01497 1.50 0.01759 1.30 0.01779 1.002022 1.20 0.02190 1.002206 1.11 0.02202 1.20 0.02190 1.20 0.02231 0.02206 1.11 0.02202 1.20 0.02190 1.20 0.02342 0.02539 0.02539 0.02539 0.02539 0.02231 0.03368 0.03873 0.04462 0.03459 0.03368 0.03873 0.04462 0.03459 0.03368 0.03873 0.04462 0.03459 0.03296 0.03368 0.03873 0.04474 0.03296 0.03368 0.03873 0.04474 0.03296 0.03368 0.03873 0.04474 0.03296 0.03368 0.03873 0.04474 0.03296 0.03368 0.03873 0.04474 0.03462 0.04484 0.04628 0.04702 0.04798	506. 26 11.59 189. 49 11.38 189. 49 11.38 186. 48 8.74 189. 11.80 4.63 199. 09 5.03 199. 77 5.06 197. 11 5.06 197. 11 5.06 197. 11 4.63 199. 11 4.63 199. 11 4.63 199. 11 4.63 199. 11 4.63 199. 11 4.63 199. 11 4.63 199. 12 4.63 199. 13 3.86 199. 14 4.63 199. 17 3.98 199. 17 3.98 190. 57 3.38 190. 57 3.35 191. 43 3.99 192. 77 2.51 193. 18 3.01 194. 18 3.01 199. 17 2.51 199. 17 2.51 199. 17 2.51 199. 17 2.51 199. 17 2.51 199. 17 2.
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	0.00050	1068.32	22.69		0.00110	495.83	13.36
	0.00067	639.20	14.15		0.00160	358.86	9.77
	0.00069 0.00031	614.66 463.77	13.82 8.51		0.00165	352.58	9.67
	0.00106	298.73	6.01		0.00214 0.00241	325.55 298.78	5.70 5.89
	0.00131	226.33	4.45		0.00252	294.45	6.18
pp + pp	0.00138 0.00139	218.49 205.22	5.99 4.65	pd + pd	0.00264 0.00265	302.38	5.51
** **	0.00140	195.77	5.61	pe pe	0.00270	290.07 291.93	5.84 5.84
261 GeV	0.00167	170.85	4.91	281 GeV	0.00282	290.72	5.64
	0.00169 0.00171	164.33 156.63	4.79 3.84		0.00366 0.00383	269.21 253.11	5.06 5.47
Normalization	0.00203	136.01	3.18	Normalization	0.00335	245.89	5.38
uncertainty	0.00215	131.15	2.81	uncertainty	0.00408	246.45	4.58
0 5 501	0.00224 0.00233	126.84 125.13	2.62 2.78	4	0.00445 0.00460	249.12 245.09	5.12 5.42
0.57%	0.00235	118.05	2.69	0.38%	0.00476	239.57	4.63
	0.00253 0.00261	112.21 115.89	2.52 2.37		0.00500	240.47	4.92
	0.00292	103.05	2.17		0.00517 0.00551	237.71 226.12	3.32 6.05
	0.00316	101.05	2.33		0.00610	217.09	4.37
	0.00343 0.00353	94.17 97.86	1.47 2.76		0.00675 0.00685	214.27 213.29	3.53
	0.00356	96.85	2.74		0.00715	208.44	3.24 5.01
	0.00359	94.88	2.17		0.00723	206.86	3.83
	0.00399 0.00402	91.73 95.68	2.34 2.39		0.00724 0.00743	205.17 208.64	4.96 4.11
	0.00425	88.20	1.79		0.00763	206.30	3.56
	0.00454 0.00475	86.88 85.24	1.45 1.70		0.00793	206.54	3.78
	0.00475	86.59	1.62		0.00801 0.00387	198.92 192.80	2.96 3.59
	0.00503	84.11	2.58		0.00930	189.98	3.41
	0.00507 0.00584	86.18 80.88	2.62 1.19	•	0.01000	180.30	4.31
	0.00504	80.81	1.63		0.01007 0.01010	176.98 179.84	2.89 4.29
	0.00659	81.33	1.63		0.0 1103	171.76	3.00
	0.00674 0.00677	77.13 81.17	2.26 1.46		0.01164	171.93	2.33
	0.00679	76.68	2.25		0.0 1228 0.0 1240	173.82 163.14	5.36 5.13
	0.00736	75.33	2.00		0.0 1250	161.50	3.24
	0.00741 0.00767	77.40 75.52	2.04 1.34		0.01409 0.01450	148.70	2.21
	0.00810	75.16	1.63		0.01494	143.70 145.88	2.61 3.05
	0.00924	75.66	1.33		0.01504	145.89	2.74
	0.00981 0.01024	73.90 72.09	1.40 1.50		0.0 1507 0.0 1694	151.33 134.92	3.12 2.49
	0.01215	70.13	1.28		0.01724	142.31	3.23
	0.01263 0.01299	71.17 67.66	1.58 1.14		0.01771	134.14	2.76
	0.01475	67.74	1.35		0.01785 0.01906	134.57 124.83	2.75 2.37
	0.01586	69.03	1.38		0.01988	118.68	2.06
	0.01622 0.01760	64.91 65.35	1.44 1.36		0.02037 0.02206	123.16 108.24	1.89
	0.01818	63.44	1,37		0.02217	113,18	2.71 2.83
•	0.01860 0.01889	61.80 65.14	1.23 1.24		0.02221	107.31	2.69
	0.01994	65.30	1.28		0.02447 0.02557	102.12 94.42	1.82 2.11
	0.02070	62.04	1.20		0.02798	91.62	2.11
	0.02142 0.02252	61.68 61.09	1.27 1.30		0.02856 0.02931	88.13 87.21	1.63
	0.02414	60.04	1.10		0.03194	78.40	1.48 1.24
•	0.02603 0.02978	58.18	1.23		0.03235	77.71	1.27
	V. UZ 3/O	\$5.42	1.57		0.03300 0.03319	79.04 79.71	2.17 2.18
					0.03483	67.97	1.39
					0.03594	67.19	1.07
					0.03885 0.03901	61.86 59.37	1.80 1.09
					0.04087	58.13	1.16
					0.04203 0.04494	52.49	1.01
					0.04494	47.16 49.56	1.44 1.49
					0.04661	45.02	0.77
					0.04735 0.04832	46.74 44.18	0.98 1.01
					0.05535	35.33	0.84
					0.05675	33.53	0.71
					0.06120 0.06258	29.83 26.87	0.66 0.67
					0.06436	25.61	0.58
					0.06879 0.070 26	21.31 21.07	0.57
					0.07026	21.07 1 5.94	0.55 0.49
							#17#

PP + PP 303 GeV Normalization uncertainty O. 50%	0. 00067 0. 00080 0. 00137 0. 00138 0. 00165 0. 00166 0. 00169 0. 00221 0. 00232 0. 00232 0. 00232 0. 00232 0. 00259 0. 00259 0. 00357 0. 00357 0. 00357 0. 00357 0. 00357 0. 00368 0. 00500 0. 00501 0. 00501 0. 00501 0. 00674 0. 00674 0. 00674 0. 00733 0. 00674 0. 00674 0. 00733 0. 00735 0. 00674 0. 00674 0. 00735 0. 00674 0. 00674 0. 00735 0. 00674 0. 00674 0. 00735 0. 00735 0. 00735 0. 00674 0. 00674 0. 00735 0. 00674 0. 00735 0. 00368 0. 00368	670.91 649.69 312.73 229.53 204.41 207.59 167.59 163.20 142.50 127.82 123.06 118.86 117.52 106.01 103.67 96.36 96.36 97.52 88.20 89.20 80 80 80 80 80 80 80 80 80 80 80 80 80	13.86 13.569 4.25 5.59 3.445 5.59 3.185 4.45 1.07 2.10 2.30 2.70 1.60 2.73 1.61 1.74 1.70 1.74 1.70 1.32 1.32 1.32 1.32 1.32 1.32 1.32 1.32		0.00107 0.00161 0.00165 0.00214 0.00241 0.00241 0.00265 0.00265 0.00271 0.00289 0.00389 0.00461 0.00477 0.00501 0.00518 0.00461 0.00477 0.00501 0.00577 0.00717 0.00717 0.00717 0.00717 0.00718 0.00744 0.00764 0.00803 0.01010 0.01010 0.01013 0.01012 0.01012 0.01012 0.01013 0.01012 0.01010 0.0	489.67 436.35 365.51 367.62 305.67 283.63 289.50 291.25 296.09 255.56 251.78 238.39 242.18 222.78 214.22 214.82 214.82 214.82 214.82 214.82 214.83 185.70 176.82 170.83 187.12 188.45 151.03 163.46 170.83 163.47 170.83 187.12 188.70 176.82 170.83 187.12 188.70 176.82 170.83 187.12 188.70 176.82 170.83 187.12 188.70 176.82 170.83 187.12 188.70 176.82 170.83 187.12 188.70 176.82 170.83 187.12 188.70 176.82 170.83 187.72 170.83 187.72 170.83 187.72 170.83 187.72 170.83 187.72 170.83 187.72 170.83 187.72 170.83 187.72 170.83 187.72 170.83 187.72 170.83	12.87 19.555 5.25 5.25 5.36 5.36 5.36 5.36 6.36 6.36 6.36 6.3
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	0.00047	1189.59	25.55		0.00083	623.92	9.80
	0.00064	659.43	14.99		0.00108	467.53	10.89
	0.00100	318.31	6.52		0.00159	339.93	14.24
	0.00125	237.35 219.13	4.69		0.00163	335.24	6.16
	0.00131 0.00133	211.03	6.10		0.00188	326.32	10.37
	0.00160	177.24	3.78 5.10		0.00224	313.28	10.55
	0.00162	164.27	4.85		0.00229	295.53	10.45
pp + pp	0.00164	172.79	4.26		0.00254	282.28	4.16
P. P. P.	0.00195	138.26	3.27	pd → pd	0.00262	294.11	4.47
398 GeV	0.00206	132.95	2.90	397 GeV	0.00290	278.01	4.88
	0.00216	133.34	2.75	397 GET	0.00300 0.00305	269.49 274.51	7.26
	0.00225	125.61	2.82		0.00328	260.60	11.54
Normalization	0.00227	124.53	2.81	Normalization		258.14	5.50 6.04
uncertainty	0.00245	121.57	2.76	uncertainty	0.00431	255.22	6.28
	0.00252	116.95	2.41	onder to mey	0.00437	251.88	6.22
0.52%	0.00283	105.91	2.26	0.36%	0.00438	239.75	3.36
	0.00307	103.98	2,41	0.00%	0.00472	239.26	3.15
	0.00343	100.37	2.85		0.00485	239.41	2.88
	0.00346	96.01	2.78		0.00534	234.34	2.75
	0.0 0349	96.61	2.22		0.00548	238.19	4.44
	0.00392	93.52	2.39	·	0.00633	222.26	2.83
	0.00414	91.91	1.86	·	0.00714	213.64	3.38
	0.00443	90.80	1.51		0.00721	212.77	4.58
	0.00463	87.03	1.74		0.00792	206.32	5.04
	0.00474	87.90	1.66		0.00909	195.55	2.44
	0.00491	86.13	2.65		0.00988	195.93	6.87
	0.00495	82.92	2.60		0.009 96	187.85	4.49
	0.00527	85.95	1.45		0.00997	191.70	6. 68
	0.00571	83.91	1.23		0.01049	180.48	2.35
	0.00 604 0.00 646	83.43 83.82	1.69		0.01237	175.39	4.45
	0.00661	79.44	1.67 2.32		0.01284	162.43 165.51	3.16
	0.00664	83.28	1.50		0.01306	165.51	5.25
	0.00665	81.49	2.35		0.01398	157.18	2.02
	0.00722	81.04	2.16		0.01505 0.01580	148.92	2.86
	0.00727	81.43	2.17		0.01592	147.99 143.27	6.34
	0.00752	79.59	1.41		0.01543	137.99	6.33 2.45
	0.00796	79.01	1.71		0.01667	139, 15	2.73
	0.00964	73.31	1.40		0.01783	130.20	2.18
	0.01008	74.96	1.55		0.01949	123.30	3.01
	0.01085	75.38	1,45		0.02191	118.09	3.42
	0.01197	71.18	1.30		0.02203	115.59	2.27
	0.01246	73.31	1.61		0.02556	104.03	3.18
	0.01281	70.31	1,21	•	0.02654	95.34	3.37
	0.01299	74.57	1.61		0.02670	97.89	3.37
	0.01455	68.99	1.36		0.03040	85.24	2.81
	0.01509	69.20	1.63		0.03057	.84.98	2.76
	0.01567	70.83	1.42		0.03127	80.03	1.47
	0.01602	65.71	1.48		0.03160	79.56	1.27
	0.01738	66.30	1.39		0.03320	75.31	2.68
	0.01797 0.01838	66,29 6 4.06	1.45		0.03546	66.34	1.20
	0.01868		1.28		0.03870	61.62	2.39
	0.01972	68.03 66.25	1,30 1.29		0.03885	57.44	1.22
	0.02047	65.65	1.27		0.03888	63.82	2.43
	0.02119	64.04	1.27 1.32		0.04167	54.07	1.03
	0.02230	60.87	1.29		0.0447B	47.73	1.85
	0.02391	60.66	1.13		0.04495	49.42	1.16
	0.02579	59.26	1,26		0.04498 0.05153	52.43	1.87
			.,		0.05175	41.73 40.83	1.74 1.20
				•	0.05857	33.85	1.20
					0.05876	33.65	1.10
					0.06602	26.39	1.00
					0.06627	26.10	0.98
					-1	-9.10	V. 30

TABLE II
Analysis Results

E _{lab}	$\rho = \frac{\text{Re } A}{\text{Im } A}$	Δρ	Δρ/Δb (GeV/c) ²	Δρ/Δσ (mb) -1
	Reaction pp + pp		(Systematic error in ppp,	± 0.009)
52	-0.153	0.012	0.038	-0.033
80	-0.096	0.010	0.036	-0.035
199	-0.034	0.009	0.025	-0.039
261	-0.009	0.009	0.023	-0.032
303	-0.011	0.008	0.024	-0.083
398	+0.012	0.009	0.022	-0.028
	Reaction pp → pd		(Systematic error in ppd'	± 0.008)
49	-0.110	0.013	0.059	-0.011
82	-0.113	0.010	0.065	-0.011
182	-0.037	0.009	0.058	-0.008
281	-0.031	0.011	0.060	-0.007
379	+0.003	0.010	0.046	-0.007
397	+0.034	0.010	0.057	-0.008

TABLE III

Reaction pd + pd Glauber Analysis Using Eq. (16)

E lab (GeV)	ρ _{pn}	Δρ _{pn}	1 _G (mb) -1	ΔI _G (mb) -1
49	-0.081	0.018	0.0281	0.0011
82	-0.127	0.013	0.0298	0.0006
182	-0.065	0.014	0.0333	0.0005
281	-0.084	0.020	0.0348	0.0006
379	-0.045	0.017	0.0362	0.0005
397	+0.021	0.017	0.0346	0.0007

TABLE IV

Reaction pd + pd Glauber Analysis Using Eq. (19)

Elab (GeV)	ρ _{pn}	Δρ _{pn}
82	-0.127	0.021
182	-0.031	0.019
281	-0.052	0.023
379	-0.003	0.021

FIGURE CAPTIONS

- Fig. 1: Schematic layout of experiment. The H₂/D₂ gas jet fires vertically (perpendicular to figure). The detectors are on a movable ion guide and at a constant distance of 7.5 m from the beam-jet intersection point.
- Fig. 2: Electronics schematic.
- Fig. 3: Typical deuteron elastic peak vs. momentum distribution.
- Fig. 4: Proton-proton differential cross section at beam energy 398 GeV. Fitted parameters: Normalization 0.3593 \pm 0.0019 $\frac{mb}{(\text{GeV/c}^2)}, p = 0.0121 \pm 0.0085, \chi^2/N_{pts} = 76/61.$
- Fig. 5: Proton-deuteron differential cross section at beam energy 379 GeV. Fitted parameters: Normalization 0.2647 \pm 0.0008 $\frac{mb}{(\text{GeV/c}^2)}, \ \rho = 0.003 \pm 0.010, \ \chi^2/N_{\text{pts}} = 96/79.$
- Fig. 6: The ratio of the real to the imaginary part of the forward p-p nuclear amplitude. The curves are I(dashed) one Pomeron formula ReA = $\frac{\pi}{2} \frac{\partial \text{ Im A}}{\partial \text{ In S}}$, II(solid) a dispersion relation calculation discussed in text. o Ref. 17

 Ref. 5. X Present experiment.
- Fig. 7: The ratio of the real to the imaginary part of the forward p-d nuclear amplitude. The solid line is the best fit $\rho(s) = -0.450 + 0.070 \text{ ln S}_{pp}$ to the proton-deuteron results. The dashed line is our fit $\rho(s) = -0.490 + 0.076 \text{ ln S}_{pp}$ to the proton-proton results. $o \text{Ref. 6.} \times \text{Present}$ experiment.
- Fig. 8: The ratio of the real to imaginary part of the forward

p-n amplitude. The error corridor shown is from the empirical fit $\rho(s) = (-0.490 \pm 0.034) + (0.076 \pm 0.006)$ In S_{pp} to the proton-proton results. o - Ref. 8. X - Present experiment.

Fig. 9: The Glauber integral, I_{G} , as a function of energy.

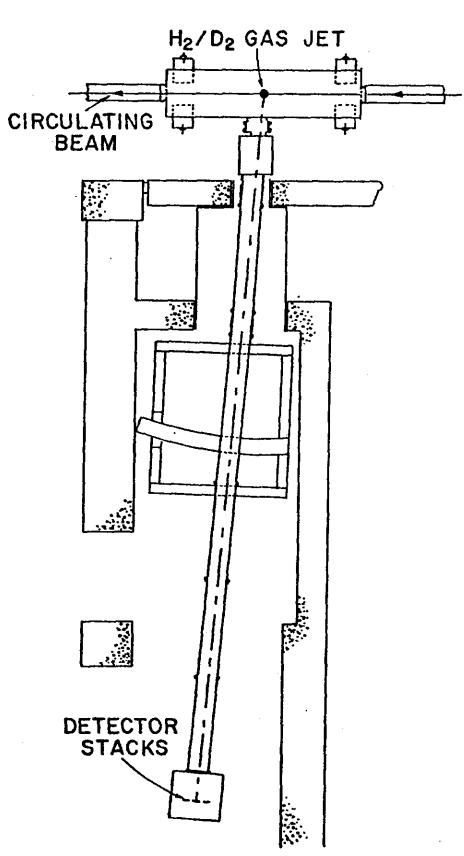
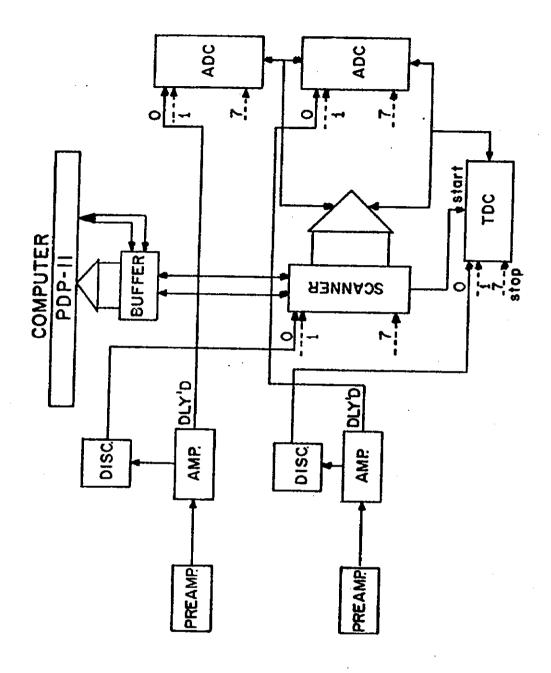
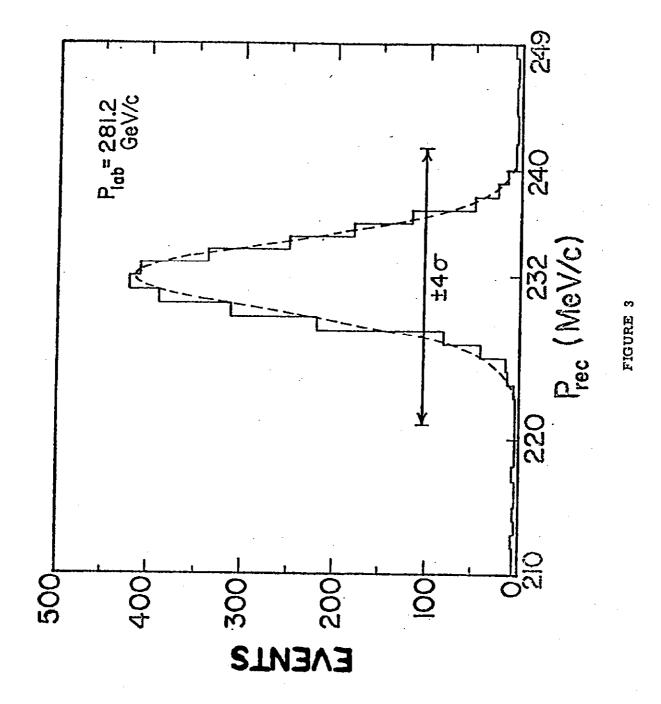


Figure 1



Pigure 2



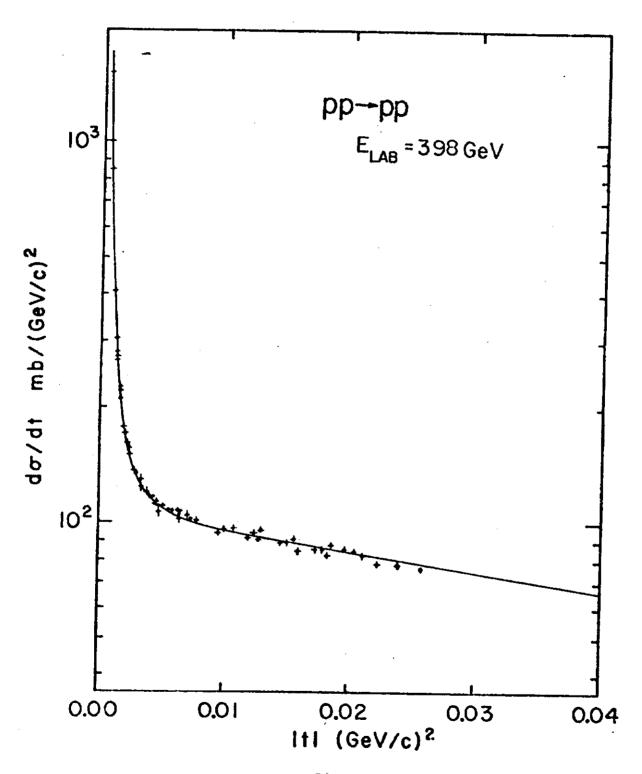


Figure 4

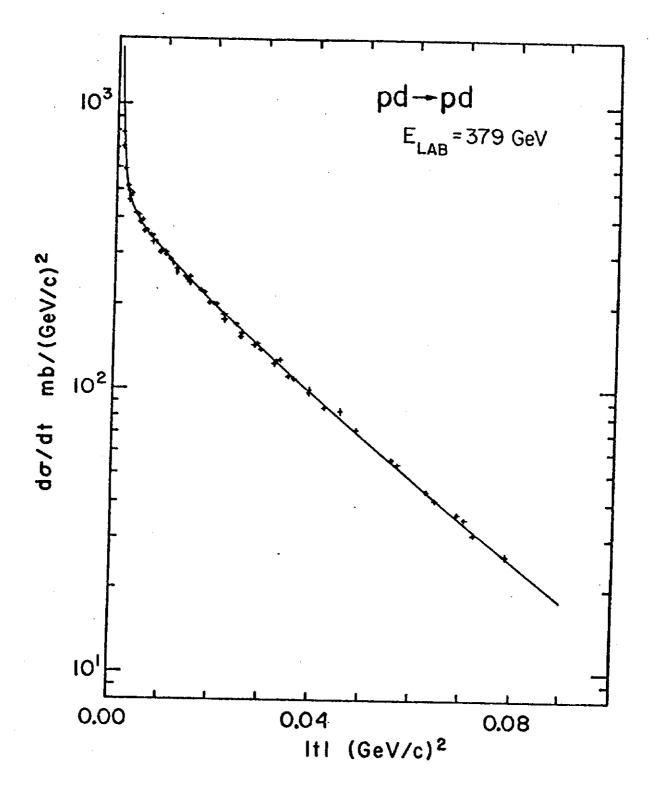


Figure 5

